

# Diffraction

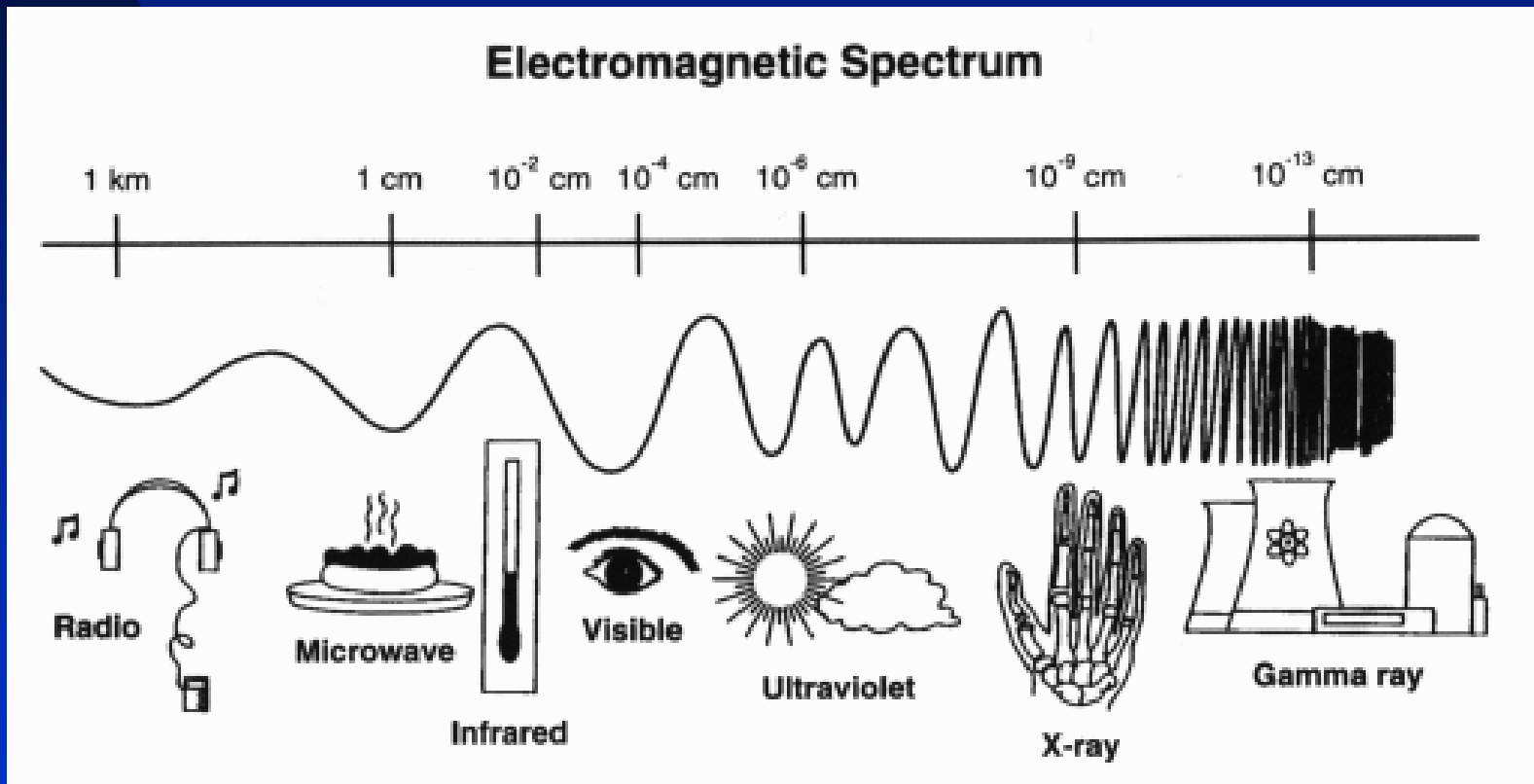
T. Ishikawa

## Part 1 Kinematical Theory

# Introduction

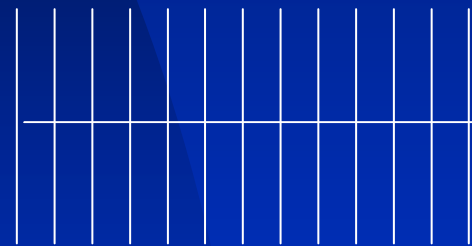
- What I intend to give in this lecture:
  - ◆ Basic Concepts of "Diffraction"
    - ★ Simpler Case: Kinematical Theory (Part 1)
    - ★ Complicated Case: Dynamical Theory (Part 2)
- The speaker has been worked on experimental dynamical diffraction for 25 years. Now, he is in charge of X-ray optics for SPring-8, world's largest 3rd generation synchrotron facility in Japan.

# X-Rays as Shorter Wavelength Electromagnetic Wave



# Scattering of x-rays by a point charge (Thomson Scattering)

$$I_0 = \sqrt{\frac{\epsilon_0}{\mu_0}} E_0^2$$



Point Charge,  
mass= $m$ ,  
charge= $e$



$$\mathbf{E}_{\text{radiation}} = \frac{1}{4\pi\epsilon_0 c^2} \frac{e(\ddot{\mathbf{x}} \times \mathbf{r}) \times \mathbf{r}}{r^3}$$

$$I = \sqrt{\frac{\epsilon_0}{\mu_0}} \left( \frac{e^2 \sin^2 \alpha}{4\pi m r c^2} E_0 \right)^2 = r_0^2 \sin^2 \alpha \left( \frac{I_0}{r^2} \right)$$

$$r_0 = \frac{e^2}{4\pi\epsilon_0 m c^2} : \text{Classical Electron Radius}$$

$$= 2.81776 \times 10^{-15} \text{ m}$$

**Lorentz Force**

$$m\ddot{\mathbf{x}} = e(\mathbf{E} + \dot{\mathbf{x}} \times \mathbf{B})$$

$$|\dot{\mathbf{x}}| \ll c \Rightarrow \ddot{\mathbf{x}} = \frac{e}{m} \mathbf{E}$$

**Electromagnetic Plane Wave**

$$\mathbf{E} = \mathbf{E}_0 \exp[i(\mathbf{K}_0 \cdot \mathbf{r} - \omega t)]$$

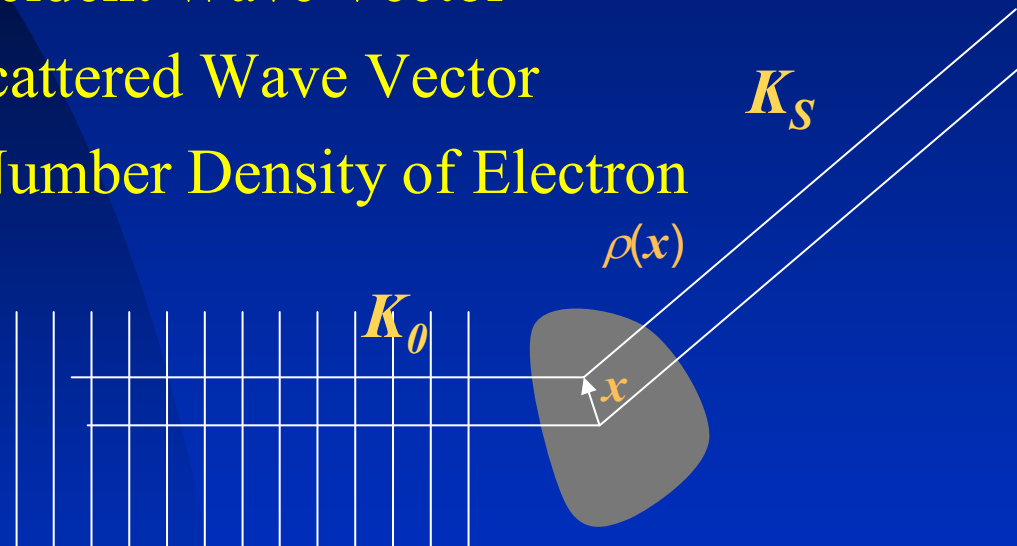
$$\mathbf{B} = \frac{\mathbf{k}}{\omega} \times \mathbf{E}_0 \exp[i(\mathbf{K}_0 \cdot \mathbf{r} - \omega t)] = \frac{\hat{\mathbf{k}}}{c} \times \mathbf{E}$$

# Scattering of x-rays by distributed charge (1/2)

$\mathbf{K}_0$  : Incident Wave Vector

$\mathbf{K}_S$  : Scattered Wave Vector

$\rho(\mathbf{x})$  : Number Density of Electron



Contribution from a volume element  $d^3\mathbf{x}$  at  $\mathbf{x}$

$$d\mathbf{E}_{\text{radiation}}^{\text{distribution}} = \mathbf{E}_{\text{radiation}}^{\text{point charge}} \rho(\mathbf{x}) \exp[-i(\mathbf{K}_S - \mathbf{K}_0) \cdot \mathbf{x}] d^3\mathbf{x}$$

# Scattering of x-rays by distributed charge (2/2)

$$\mathbf{E}_{\text{radiation}}^{\text{distribution}} = \mathbf{E}_{\text{radiation}}^{\text{point charge}} \iiint \rho(\mathbf{x}) \exp[-i(\mathbf{K}_s - \mathbf{K}_o) \cdot \mathbf{x}] d^3 \mathbf{x}$$

3D Fourier Transform of Charge Density

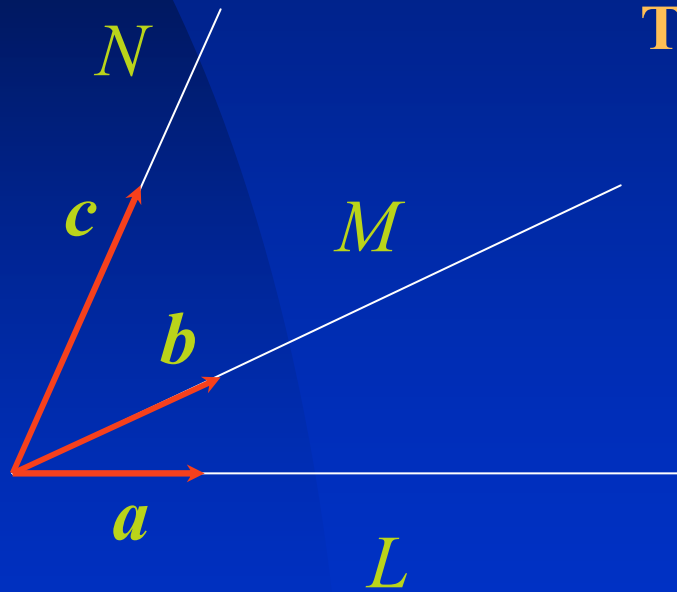
Scattered Intensity

$$\begin{aligned} I &= I^{\text{single}} \left| \iiint \rho(\mathbf{x}) \exp[-i(\mathbf{K}_s - \mathbf{K}_o) \cdot \mathbf{x}] d^3 \mathbf{x} \right|^2 \\ &= \frac{r_o^2 \sin^2 \alpha}{r^2} I_o \left| \iiint \rho(\mathbf{x}) \exp[-i(\mathbf{K}_s - \mathbf{K}_o) \cdot \mathbf{x}] d^3 \mathbf{x} \right|^2 \end{aligned}$$

# Electronic Charge Distribution in Crystal

Crystal = 3D Regular Stacking of Molecules

Translation Symmetry



$$\rho(\mathbf{x}) = \rho(\mathbf{x} + l\mathbf{a} + m\mathbf{b} + n\mathbf{c})$$

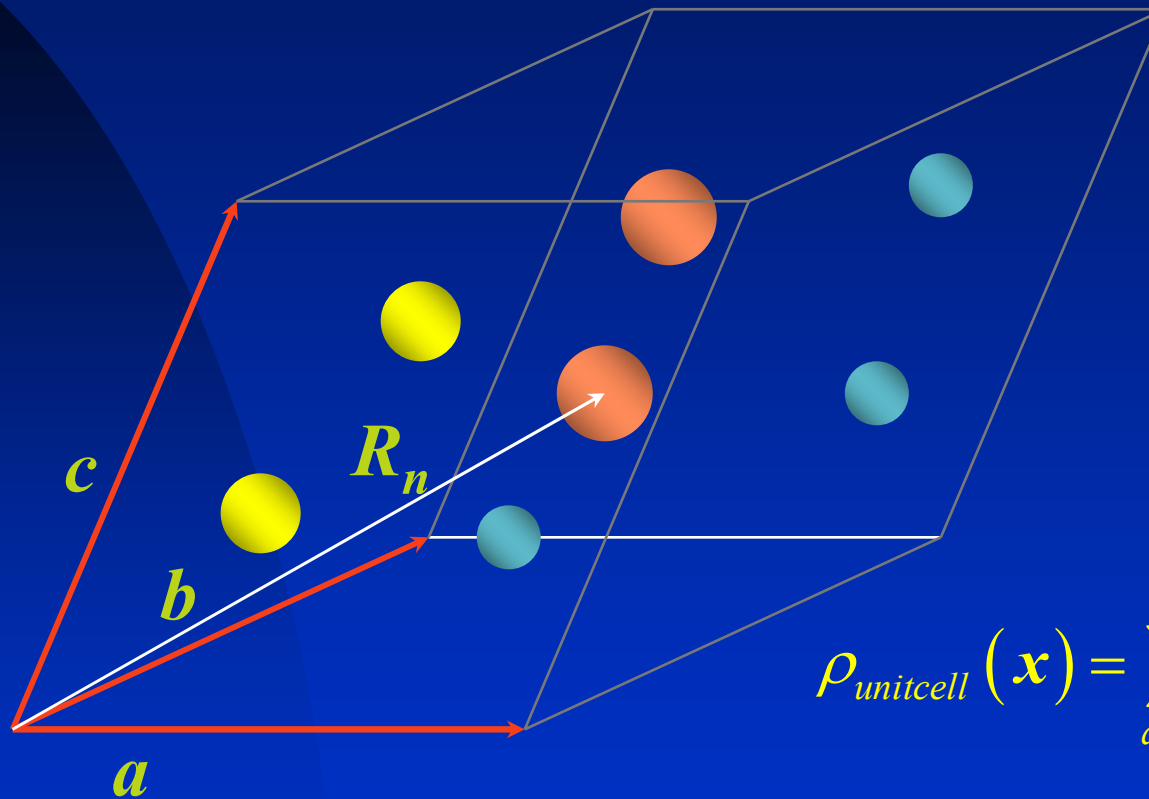
$l, m, n$ : integers

# Fourier Transform of the Electronic Charge Distribution in Crystal

$$\begin{aligned} \int_{\text{crystal}} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} &= \sum_{l=0}^{L-1} \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} \int_{\text{unitcell}} \rho(\mathbf{x} + l\mathbf{a} + m\mathbf{b} + n\mathbf{c}) e^{-i\mathbf{K}\cdot(\mathbf{x} + l\mathbf{a} + m\mathbf{b} + n\mathbf{c})} d^3\mathbf{x} \\ &= \sum_{l=0}^{L-1} \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} e^{-i\mathbf{K}\cdot(l\mathbf{a} + m\mathbf{b} + n\mathbf{c})} \int_{\text{unitcell}} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} \\ &= \left( \frac{1 - e^{-iL\mathbf{K}\cdot\mathbf{a}}}{1 - e^{-i\mathbf{K}\cdot\mathbf{a}}} \right) \left( \frac{1 - e^{-iM\mathbf{K}\cdot\mathbf{b}}}{1 - e^{-i\mathbf{K}\cdot\mathbf{b}}} \right) \left( \frac{1 - e^{-iN\mathbf{K}\cdot\mathbf{c}}}{1 - e^{-i\mathbf{K}\cdot\mathbf{c}}} \right) \int_{\text{unitcell}} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} \end{aligned}$$



# Charge Density in Unit Cell



$$\rho_{\text{unitcell}}(\mathbf{x}) = \sum_{\text{atom}} \rho_{\text{atom}}^{(a)}(\mathbf{x} - \mathbf{R}_n)$$

$$\int_{\text{unitcell}} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} = \sum_{\text{atom}} \int_{\text{atom}} \rho_{\text{atom}}^{(a)}(\mathbf{x} - \mathbf{R}_n) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} = \sum_{\text{atom}} e^{-i\mathbf{K}\cdot\mathbf{R}_n} \int_{\text{atom}} \rho_{\text{atom}}^{(a)}(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x}$$

# Atomic Scattering Factor, Structure Factor

## Atomic Scattering Factor

$$f^{(a)}(\mathbf{K}) = \int_{atom} \rho_{atom}^{(a)}(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x}$$

## Structure Factor

$$F(\mathbf{K}) = \int_{unitcell} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} = \sum_{atom} e^{-i\mathbf{K}\cdot\mathbf{R}_n} f^{(a)}(\mathbf{K})$$

$$\begin{aligned} \int_{crystal} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} &= \left( \frac{1 - e^{-iLK\cdot a}}{1 - e^{-iK\cdot a}} \right) \left( \frac{1 - e^{-iMK\cdot b}}{1 - e^{-iK\cdot b}} \right) \left( \frac{1 - e^{-iNK\cdot c}}{1 - e^{-iK\cdot c}} \right) \int_{unitcell} \rho(\mathbf{x}) e^{-i\mathbf{K}\cdot\mathbf{x}} d^3\mathbf{x} \\ &= \left( \frac{1 - e^{-iLK\cdot a}}{1 - e^{-iK\cdot a}} \right) \left( \frac{1 - e^{-iMK\cdot b}}{1 - e^{-iK\cdot b}} \right) \left( \frac{1 - e^{-iNK\cdot c}}{1 - e^{-iK\cdot c}} \right) F(\mathbf{K}) \end{aligned}$$

# Laue Function (1/3)

$$\frac{1 - e^{-iLK \cdot a}}{1 - e^{-iK \cdot a}} = \frac{e^{-i\frac{LK \cdot a}{2}} \left( e^{+i\frac{LK \cdot a}{2}} - e^{-i\frac{LK \cdot a}{2}} \right)}{e^{-i\frac{K \cdot a}{2}} \left( e^{+i\frac{K \cdot a}{2}} - e^{-i\frac{K \cdot a}{2}} \right)} = \frac{e^{-i\frac{LK \cdot a}{2}} \sin \frac{LK \cdot a}{2}}{e^{-i\frac{K \cdot a}{2}} \sin \frac{K \cdot a}{2}}$$



$$I_{crystal} = \frac{r_o^2 \sin^2 \alpha}{r^2} I_o |F(K)|^2 \left( \frac{\sin \frac{LK \cdot a}{2}}{\sin \frac{K \cdot a}{2}} \right)^2 \left( \frac{\sin \frac{MK \cdot b}{2}}{\sin \frac{K \cdot b}{2}} \right)^2 \left( \frac{\sin \frac{NK \cdot c}{2}}{\sin \frac{K \cdot c}{2}} \right)^2$$

# Laue Function (2/3)

$$\mathbf{K} \cdot \mathbf{a} \rightarrow 0 \Rightarrow \left( \frac{\sin \frac{L\mathbf{K} \cdot \mathbf{a}}{2}}{\sin \frac{\mathbf{K} \cdot \mathbf{a}}{2}} \right)^2 \rightarrow L^2$$

$$\left( \frac{\sin \frac{L\mathbf{K} \cdot \mathbf{a}}{2}}{\sin \frac{\mathbf{K} \cdot \mathbf{a}}{2}} \right)^2 : \text{periodic function}$$

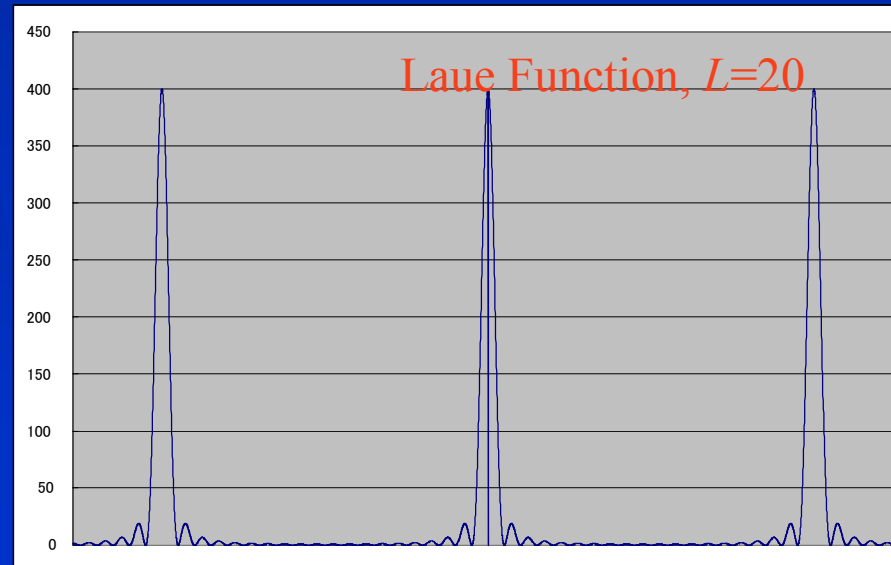
$$\mathbf{K} \cdot \mathbf{a} \rightarrow 2h\pi \Rightarrow \left( \frac{\sin \frac{L\mathbf{K} \cdot \mathbf{a}}{2}}{\sin \frac{\mathbf{K} \cdot \mathbf{a}}{2}} \right)^2 \rightarrow L^2 ; h \text{ integer}$$

Similarly,

$$\mathbf{K} \cdot \mathbf{b} \rightarrow 2k\pi, k \text{ integer} \Rightarrow \left( \frac{\sin \frac{M\mathbf{K} \cdot \mathbf{b}}{2}}{\sin \frac{\mathbf{K} \cdot \mathbf{b}}{2}} \right)^2 \rightarrow M^2$$

and

$$\mathbf{K} \cdot \mathbf{c} \rightarrow 2l\pi, l \text{ integer} \Rightarrow \left( \frac{\sin \frac{N\mathbf{K} \cdot \mathbf{c}}{2}}{\sin \frac{\mathbf{K} \cdot \mathbf{c}}{2}} \right)^2 \rightarrow N^2$$



## Laue Function (3/3)

- Scattering from a crystal is appreciable only when

$$\mathbf{K} \cdot \mathbf{a} = 2\pi h$$

$$\mathbf{K} \cdot \mathbf{b} = 2\pi k \quad h, k, l; \text{ integer}$$

$$\mathbf{K} \cdot \mathbf{c} = 2\pi l$$

- Then,

$$I_{crystal} = \frac{r_o^2 \sin^2 \alpha}{r^2} I_o (LMN)^2 |F(\mathbf{K})|^2$$

# Reciprocal Lattice (1/3)

Scattering from a crystal is appreciable only when

$$\mathbf{K} \cdot \mathbf{a} = 2\pi h$$

$$\mathbf{K} \cdot \mathbf{b} = 2\pi k$$

$$\mathbf{K} \cdot \mathbf{c} = 2\pi l$$

$h, k, l$  : integer

## Base Vectors of Reciprocal Lattice

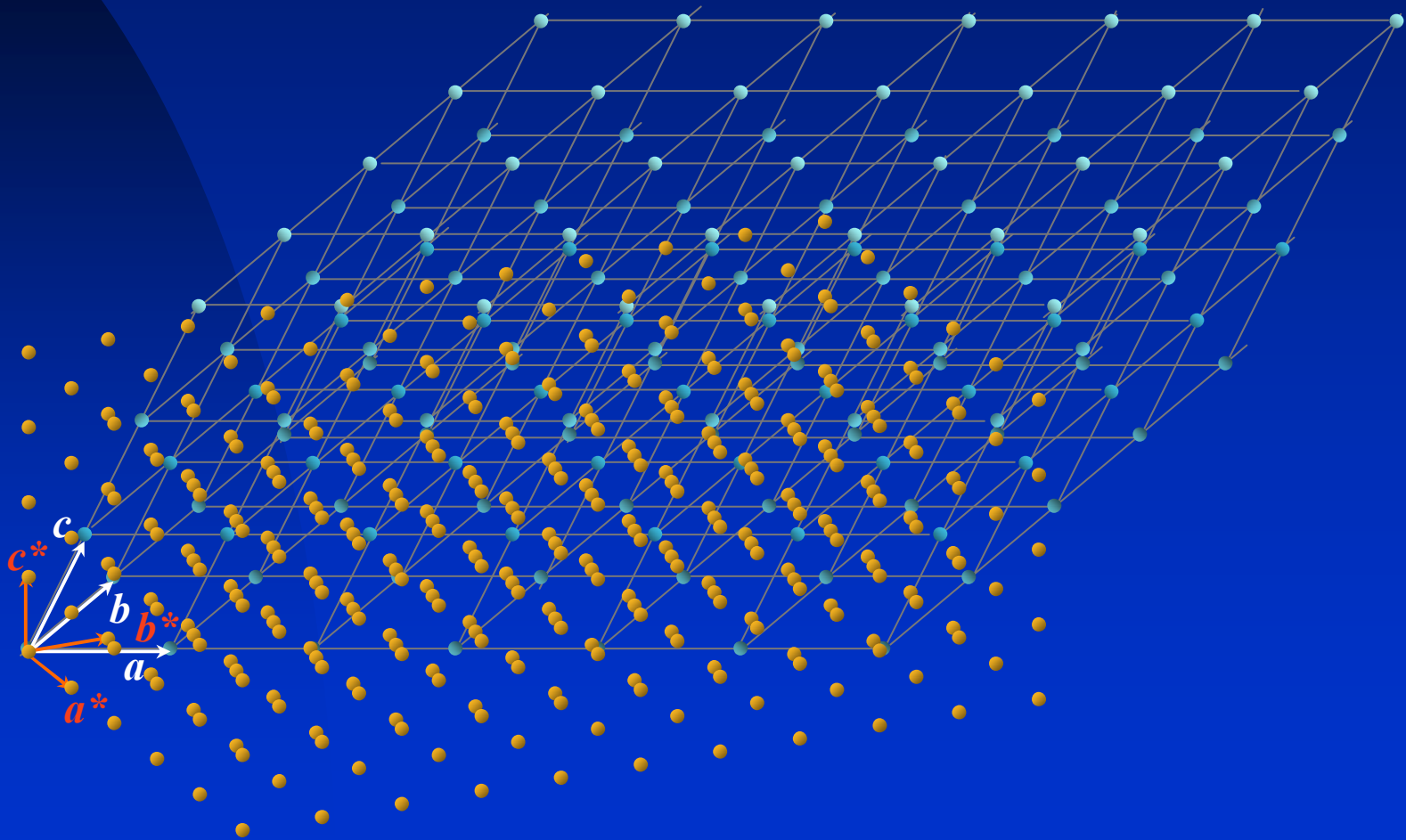
$$\mathbf{a}^* = \frac{2\pi(\mathbf{b} \times \mathbf{c})}{\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c})}, \mathbf{b}^* = \frac{2\pi(\mathbf{c} \times \mathbf{a})}{\mathbf{b} \cdot (\mathbf{c} \times \mathbf{a})}, \mathbf{c}^* = \frac{2\pi(\mathbf{a} \times \mathbf{b})}{\mathbf{c} \cdot (\mathbf{a} \times \mathbf{b})}$$

$$\mathbf{a}^* \cdot \mathbf{a} = 2\pi, \mathbf{a}^* \cdot \mathbf{b} = 0, \mathbf{a}^* \cdot \mathbf{c} = 0$$

$$\mathbf{b}^* \cdot \mathbf{b} = 2\pi, \mathbf{b}^* \cdot \mathbf{a} = 0, \mathbf{b}^* \cdot \mathbf{c} = 0$$

$$\mathbf{c}^* \cdot \mathbf{c} = 2\pi, \mathbf{c}^* \cdot \mathbf{a} = 0, \mathbf{c}^* \cdot \mathbf{b} = 0$$

# Reciprocal Lattice (2/3)



## Reciprocal Lattice (3/3)

- Reciprocal Lattice Vector

$$\mathbf{g} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*$$

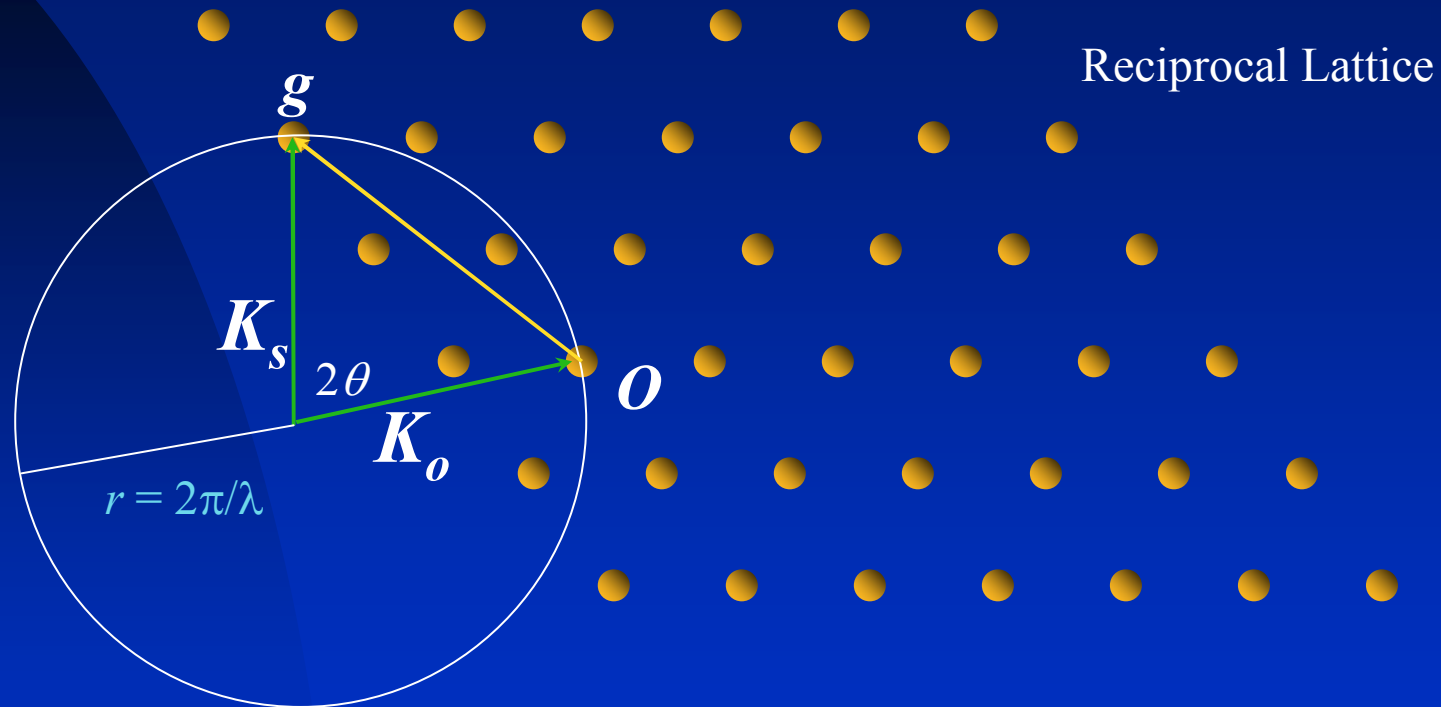
$h, k, l$ ; integer

$$\mathbf{K} = \mathbf{g} \Rightarrow \mathbf{K} \cdot \mathbf{a} = 2\pi h, \mathbf{K} \cdot \mathbf{b} = 2\pi k, \mathbf{K} \cdot \mathbf{c} = 2\pi l$$

When the scattering vector,  $\mathbf{K} = \mathbf{K}_s - \mathbf{K}_o$ , corresponds to a reciprocal lattice vector, strong diffraction may be observed (necessary condition, but not a sufficient condition).



# Ewald Sphere



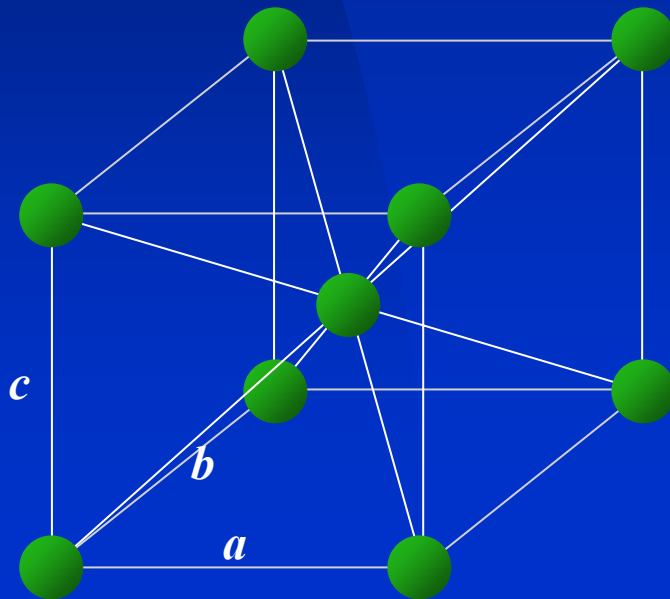
$$|\mathbf{g}| = \frac{2\pi}{d_{hkl}}, |\mathbf{K}_o| = |\mathbf{K}_s| = \frac{2\pi}{\lambda} \Rightarrow \frac{2\pi}{\lambda} \sin \theta = \frac{\pi}{d_{hkl}} \quad \text{Bragg Condition of Diffraction}$$
$$\Rightarrow \lambda = 2d_{hkl} \sin \theta$$

# Forbidden Reflection (1/2)

$\mathbf{K}=\mathbf{g}$  is a necessary condition for observing diffraction, but not a sufficient condition...

If  $F(\mathbf{K})=0$ , then  $I_{crystal} = 0$  even when  $\mathbf{K}=\mathbf{g}$ .

Example 1, Body Center Cubic Lattice (bcc)



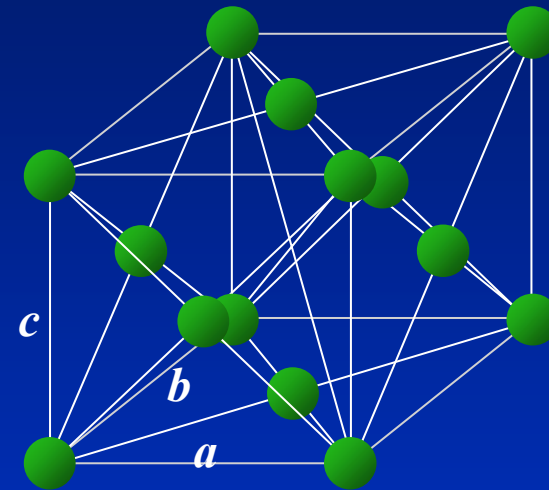
$$F(\mathbf{g}) = f^a(\mathbf{g}) \left\{ \exp(i\mathbf{g} \cdot \mathbf{0}) + \exp\left(i\mathbf{g} \cdot \left(\frac{1}{2}\mathbf{a} + \frac{1}{2}\mathbf{b} + \frac{1}{2}\mathbf{c}\right)\right) \right\}$$
$$= f^a(\mathbf{g}) \left\{ 1 + \exp(i\pi(h+k+l)) \right\}$$

$$h+k+l = \text{odd} \Rightarrow F(\mathbf{g}) = 0 \text{ (Forbidden Reflection)}$$

$$h+k+l = \text{even} \Rightarrow F(\mathbf{g}) = 2f^a(\mathbf{g})$$

# Forbidden Reflection (2/2)

Example 2, Face Center Cubic Lattice (fcc)



$$F(\mathbf{g}) = f^a(\mathbf{g}) \left\{ 1 + \exp(i\pi(h+k)) + \exp(i\pi(k+l)) + \exp(i\pi(l+h)) \right\}$$

$$h, k, l \text{ all odd or all even} \Rightarrow F(\mathbf{g}) = 4f^a(\mathbf{g})$$

$$\text{Otherwise} \Rightarrow F(\mathbf{g}) = 0 \text{ (Forbidden Reflection)}$$

# Special Topics

**What we can measure in diffraction/scattering experiment  
= Intensity**

**All phase information is lost!**

Non-Crystalline Charge Distribution:

$$I = \frac{r_o^2 \sin^2 \alpha}{r^2} I_o \left( \iiint \rho(\mathbf{x}) \exp[-i\mathbf{K} \cdot \mathbf{x}] d^3 \mathbf{x} \right)^2 = \sqrt{\frac{\epsilon_o}{\mu_o}} \mathbf{E} \cdot \mathbf{E}^*$$

$\mathbf{E}^*$  : complex conjugate of  $\mathbf{E}$

If

$$\begin{aligned} E &= \frac{r_o \sin \alpha}{r} \sqrt{I_o \sqrt{\frac{\mu_o}{\epsilon_o}}} \iiint \rho(\mathbf{x}) \exp[-i\mathbf{K} \cdot \mathbf{x}] d^3 \mathbf{x} \\ &= \sqrt{I} \sqrt{\frac{\mu_o}{\epsilon_o}} \exp(i\varphi) \end{aligned}$$

is obtained, we can calculate  $\rho(\mathbf{x})$  by Fourier inversion.

# Iterative Phase Retrieval

(Jianwei Miao & David Sayre)

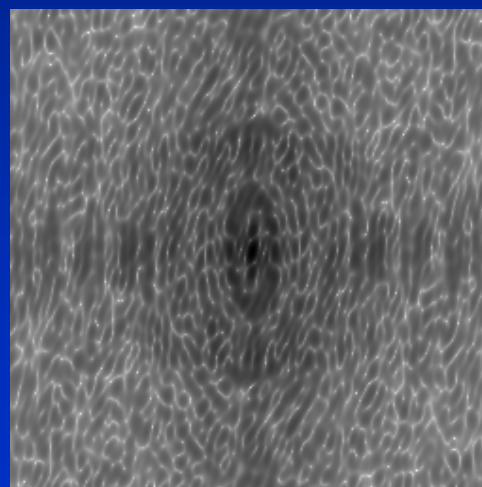
**X-ray intensity data: Phase Information is Lost!**

**Scattered pattern in Far Field with Coherent Illumination, Phase can be retrieved.**

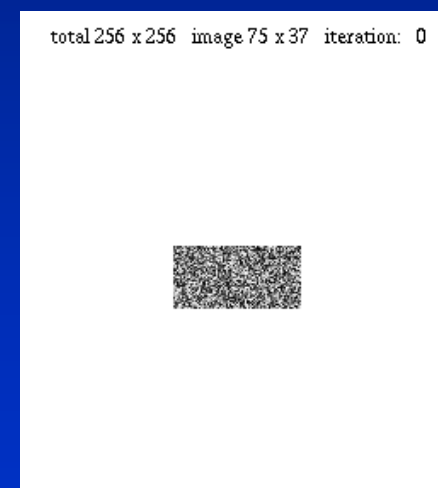
**Phase Retrieval → Iterative Algorithm developed by Gerchberg & Saxon, followed by the improvement by Fienup (Opt. Lett. 3 (1978) 27.)**



**Real Space Image**



**Scattered Intensity**

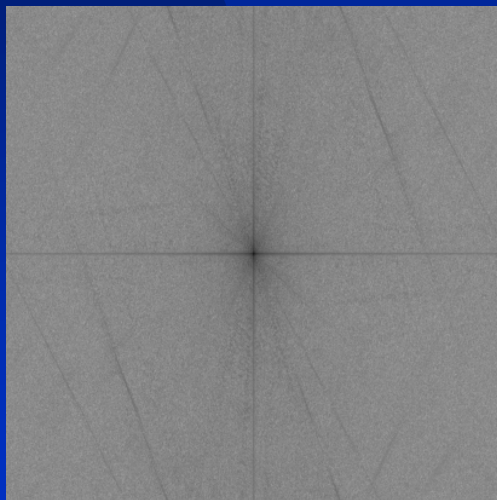


**Phase Retrieval**

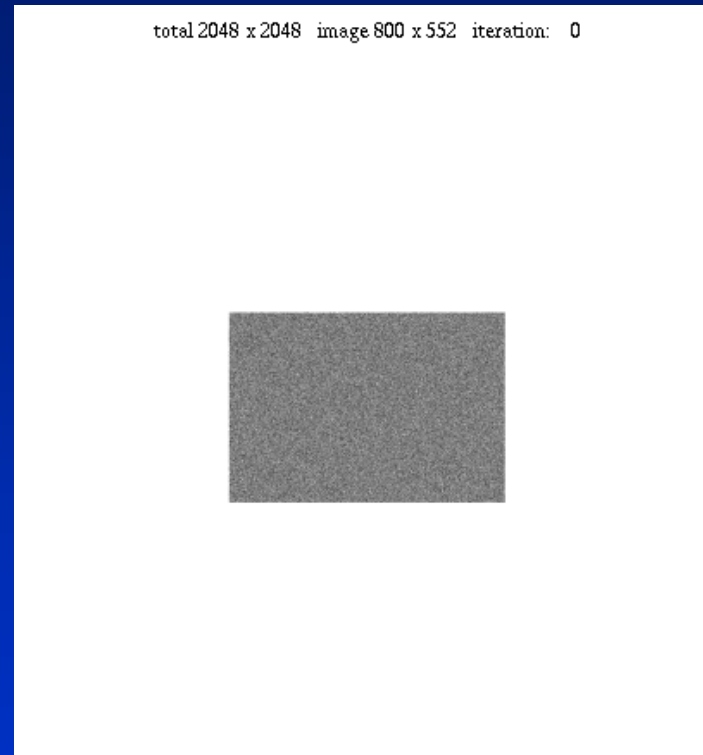
# Reconstruction of Complex Real Space Images



**Real Space**



**Scattered Intensity**



**Phase Retrieval**



**Original Image**

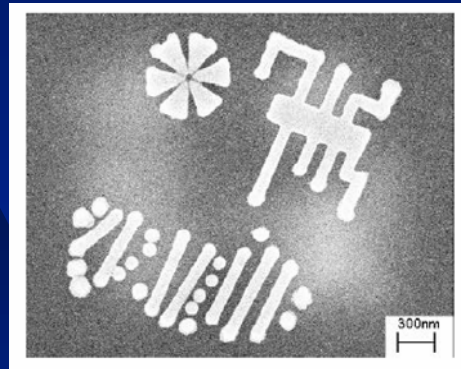
**Reconstructed Image  
After 5000 iteration**



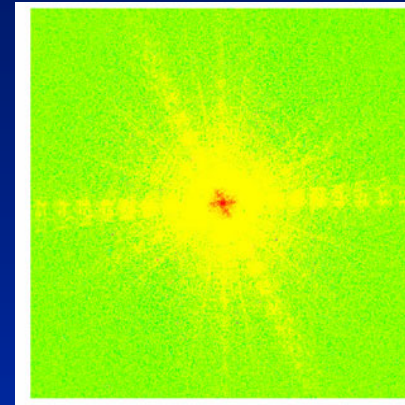
# 3D Diffraction Microscopy

Miao *et al.* PRL (2002)

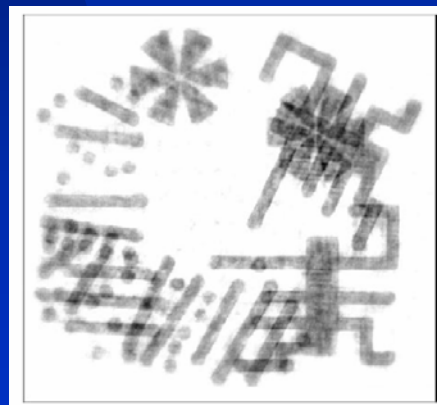
## Two Layer Ni Pattern



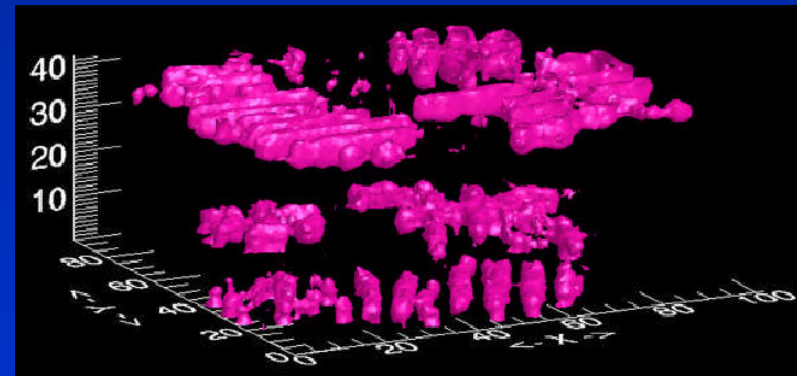
SEM image of Ni pattern on SiN



Coherent Scattering Pattern



2D Reconstructed Image  
( $<10$  nm resolution)



3D Reconstructed Image ( $\sim 50$  nm resolution)